

## Chapter 6

### Application of the Distributions $e^{\alpha t} \square^k \delta$ to the Telegraph Equation

In this chapter, we introduce a new technique for solving the telegraph equation by using the distribution  $e^{\alpha t} \square^k \delta$ .

#### 6.1 Main Results

Consider the Telegraph equation

$$T_1^k u(x, t) = \delta(x, t) \quad (6.1)$$

and

$$T_2^k u(x, t) = \delta(x, t) \quad (6.2)$$

where  $T_1^k$  and  $T_2^k$  are defined by (1.33) and (1.34), respectively,  $(x, t) \in \mathbb{R}^n \times (0, \infty)$ .

Changing the independent variables from  $x_1$  to  $t$ , where  $t$  is the time and putting  $p = 1$  and  $q = n$ , then the  $n$ -dimensional to be  $n + 1$ -dimensional and choosing  $\alpha_2 = \alpha_3 \dots = \alpha_n = 0$  of the distribution  $e^{\alpha t} \square^k \delta$ , we obtain

$$e^{\alpha t} \square^k \delta = e^{\alpha_1 t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right)^k \delta(x, t), \quad (6.3)$$

where  $\Delta$  is the Laplace operator defined by (1.2).

Next, we recalling the Legendre's duplication of  $\Gamma(z)$ ,

$$\Gamma(2z) = 2^{2z-1} \pi^{-1/2} \Gamma(z) \Gamma(z + \frac{1}{2})$$

then the function (2.13) may reduces to the function

$$M_\gamma(v) = \begin{cases} \frac{v^{(\gamma-n-1)/2}}{H_{n+1}(\gamma)} & \text{if } t \in \Gamma_+, \\ 0 & \text{if } t \notin \Gamma_+ \end{cases} \quad (6.4)$$

where  $v = t^2 - x_1^2 - \dots - x_n^2$ ,  $t$  is the time and

$$H_{n+1}(\gamma) = \pi^{(n-1)/2} 2^{\gamma-1} \Gamma\left(\frac{\gamma-n+1}{2}\right) \Gamma\left(\frac{\gamma}{2}\right),$$

the function  $M_\gamma(s)$  is called the hyperbolic kernel of Marcel Riesz.

By the property of  $\delta(x, t) = \delta(x)\delta(t)$  and using the Proposition 2.1.14 of  $e^{\alpha t}\delta^{(k)} = \left(\frac{\partial}{\partial t} - \alpha\right)^k \delta$  and  $e^{\alpha_1 t}\delta = \delta$  we can express (6.3) for  $k = 1$  as

$$\begin{aligned} e^{\alpha_1 t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right) \delta(x, t) &= e^{\alpha_1 t} \frac{\partial^2}{\partial t^2} \delta - e^{\alpha_1 t} \Delta \delta \\ &= \left( \frac{\partial}{\partial t} - \alpha_1 \right)^2 \delta - \Delta e^{\alpha_1 t} \delta \\ &= \left( \frac{\partial^2}{\partial t^2} - 2\alpha_1 \frac{\partial}{\partial t} + \alpha_1^2 - \Delta \right) \delta. \end{aligned}$$

Choosing  $\alpha_1 = -\beta$ , it follows that

$$e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right) \delta(x, t) = \left( \frac{\partial^2}{\partial t^2} + 2\beta \frac{\partial}{\partial t} + \beta^2 - \Delta \right) \delta(x, t) = T_1 \delta(x, t). \quad (6.5)$$

Now, convolution both sides of (6.5) by the distribution  $e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right) \delta(x, t)$ ,  $k$ -times we obtain

$$e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right) \delta(x, t) * \dots * e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right) \delta(x, t) = T_1 \delta * \dots * T_1 \delta. \quad (6.6)$$

Then we have

$$e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right)^k \delta(x, t) = T_1^k \delta. \quad (6.7)$$

Solving the Telegraph equation (6.1) after substituting (6.7) in to (6.1) and using the Lemma 2.4.9, thus

$$T_1^k u(x, t) = e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} - \Delta \right)^k \delta(x, t) * u(x, t) = \delta(x, t), \quad (6.8)$$

convolving both side by  $e^{-\beta t} M_{2k}(v)$  of (6.8), we obtain

$$u(x, t) = e^{-\beta t} M_{2k}(v), \quad (6.9)$$

the elementary solution of the Telegraph equation of the (6.1).

We next consider the distribution  $e^{\alpha t} \square^k \delta$  with  $k = 1$ ,  $\alpha_1 = -\beta$ ,  $p = n + 1$ ,  $q = 0$  and  $x_1 = t$ . then we obtain

$$\begin{aligned} e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} + \Delta \right) \delta(x, t) &= e^{-\beta t} \frac{\partial^2}{\partial t^2} \delta + e^{-\beta t} \Delta \delta \\ &= \left( \frac{\partial}{\partial t} + \beta \right)^2 \delta + \Delta e^{-\beta t} \delta \\ &= \left( \frac{\partial^2}{\partial t^2} + 2\beta \frac{\partial}{\partial t} + \beta^2 + \Delta \right) \delta. \end{aligned}$$

Similarly, convolving the distribution  $e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} + \Delta \right) \delta(x, t)$ ,  $k$ -times we obtain the operator  $T_2^k$  as

$$e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} + \Delta \right)^k \delta(x, t) = T_2^k \delta, \quad (6.10)$$

where  $T_2^k$  is defined by (1.34),  $(x, t) \in \mathbb{R}^n \times (0, \infty)$ . By Legendre's duplication of  $\Gamma(z)$  with  $p = n + 1$  and  $q = 0$ , the hyperbolic kernel reduce to the elliptic kernel of Marcel Riesz denoted by  $R_\lambda(s) = s^{(\lambda-n-1)/2} / W_{n+1}(\lambda)$ , where  $s^2 = t^2 + x_1^2 + x_2^2 + \dots + x_n^2$  and  $W_{n+1}(\lambda) = \pi^{n+1/2} 2^\lambda \Gamma(\lambda/2) / \Gamma((n+1-\lambda)/2)$ . Solving the Telegraph equation (6.2) after substituting (6.10) in to (6.2) and using the Lemma 2.4.2, see also [11], we have

$$T_2^k u(x, t) = e^{-\beta t} \left( \frac{\partial^2}{\partial t^2} + \Delta \right)^k \delta(x, t) * u(x, t) = \delta(x, t), \quad (6.11)$$

convolving both side by  $e^{-\beta t} (-1)^k R_{2k}(s)$  of (6.11), we obtain

$$u(x, t) = e^{-\beta t} (-1)^k R_{2k}(s), \quad (6.12)$$

the elementary solution of the Telegraph equation of the (6.2) as required.

## 6.2 Example

In this section, we want to show an example about application of the distributions to the telegraph equation.

**Example 6.2.1** We now to find the elementary solution of the equation

$$\frac{\partial^2}{\partial t^2} u(x, t) + 2\beta \frac{\partial}{\partial t} u(x, t) + \beta^2 u(x, t) - \frac{\partial^2}{\partial x^2} u(x, t) = 0 \quad (6.13)$$

with the initial conditions

$$u(x, 0) = 0, \quad \frac{\partial}{\partial t} u(x, 0) = \delta(x) \quad (6.14)$$

where  $\delta(x)$  is the Dirac-delta function.

Now, we consider the equation

$$\frac{\partial^2}{\partial t^2} u(x, t) + 2\beta \frac{\partial}{\partial t} u(x, t) + \beta^2 u(x, t) - \frac{\partial^2}{\partial x^2} u(x, t) = \delta(x) \delta(t). \quad (6.15)$$

Taking the Fourier transform to both side of the equation (6.15)

$$\frac{\partial^2}{\partial t^2} \hat{u}(\xi, t) + 2\beta \frac{\partial}{\partial t} \hat{u}(\xi, t) + \beta^2 \hat{u}(\xi, t) + \xi^2 \hat{u}(\xi, t) = \delta(t). \quad (6.16)$$

Now consider the equation

$$\frac{\partial^2}{\partial t^2} \hat{u}(\xi, t) + 2\beta \frac{\partial}{\partial t} \hat{u}(\xi, t) + \beta^2 \hat{u}(\xi, t) + \xi^2 \hat{u}(\xi, t) = 0 \quad (6.17)$$

with the initial conditions

$$\hat{u}(\xi, 0) = 0, \quad \frac{\partial}{\partial t} \hat{u}(\xi, 0) = 1.$$

Then we obtain  $\hat{\phi}(\xi, t)$  as a solution of equation (6.17) where

$$\hat{\phi}(\xi, t) = \frac{e^{-\beta t}}{2\xi i} (\exp[i\xi t] - \exp[-i\xi t]).$$

Thus equation (6.16) has solution

$$\hat{u}(\xi, t) = H(t) \hat{\phi}(\xi, t)$$

where  $H(t)$  a Heaviside function. Since the right hand side of the equation is tempered distribution, but they are not locally integrable. Thus we can not find the inverse Fourier transform  $u(x, t)$ . So, we define

$$\hat{\phi}^\epsilon(\xi, t) = e^{-\epsilon t} \hat{\phi}(\xi, t)$$

and

$$\hat{\phi}^\epsilon(\xi, t) \rightarrow \hat{\phi}(\xi, t) \quad \text{uniformly as} \quad \epsilon \rightarrow 0.$$

Thus

$$\hat{\phi}^\epsilon(\xi, t) = \frac{e^{-\beta t}}{2i} \left[ \frac{e^{-(\epsilon - it)\xi} - e^{-(\epsilon + it)\xi}}{\xi} \right]$$

is locally integrable and then we can find the inverse Fourier transform  $\phi(x, t)$ .

Now

$$\hat{\phi}^\epsilon(\xi, t) = \frac{e^{-\beta t}}{2i} \int_{\epsilon-it}^{\epsilon+it} e^{-s\xi} ds.$$

Applying the inverse Fourier transform, we obtain

$$\begin{aligned} \phi^\epsilon(x, t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\xi x} \hat{\phi}^\epsilon(\xi, t) d\xi \\ &= \frac{e^{-\beta t}}{4\pi i} \int_{\epsilon-it}^{\epsilon+it} \int_{-\infty}^{\infty} e^{i\xi x} e^{-s\xi} ds d\xi. \end{aligned}$$

By directly computation, we have

$$\int_{-\infty}^{\infty} e^{i\xi x} e^{-s\xi} d\xi = \frac{2s}{s^2 + x^2}.$$

Thus

$$\begin{aligned} \phi^\epsilon(x, t) &= \frac{e^{-\beta t}}{4\pi i} \int_{\epsilon-it}^{\epsilon+it} \frac{2s}{s^2 + x^2} ds \\ &= \frac{e^{-\beta t}}{4\pi i} (\ln[(\epsilon + it)^2 + x^2] - \ln[(\epsilon - it)^2 + x^2]). \end{aligned}$$

Let  $\epsilon \rightarrow 0$ , we obtain

$$\phi(x, t) = \frac{e^{-\beta t}}{4\pi i} (\ln[(x^2 - t^2) + 0i] - \ln[(x^2 - t^2) - 0i]).$$

Since, we have

$$\ln[(x^2 - t^2) + 0i] - \ln[(x^2 - t^2) - 0i] = 2\pi i H(t^2 - x^2).$$

Thus, we obtain

$$\phi(x, t) = \frac{e^{-\beta t}}{2} H(t^2 - x^2).$$

Thus equation (6.13) has

$$E(x, t) = H(t) \phi(x, t) = \frac{e^{-\beta t}}{2} H(t^2 - x^2) H(t)$$

as elementary solution.

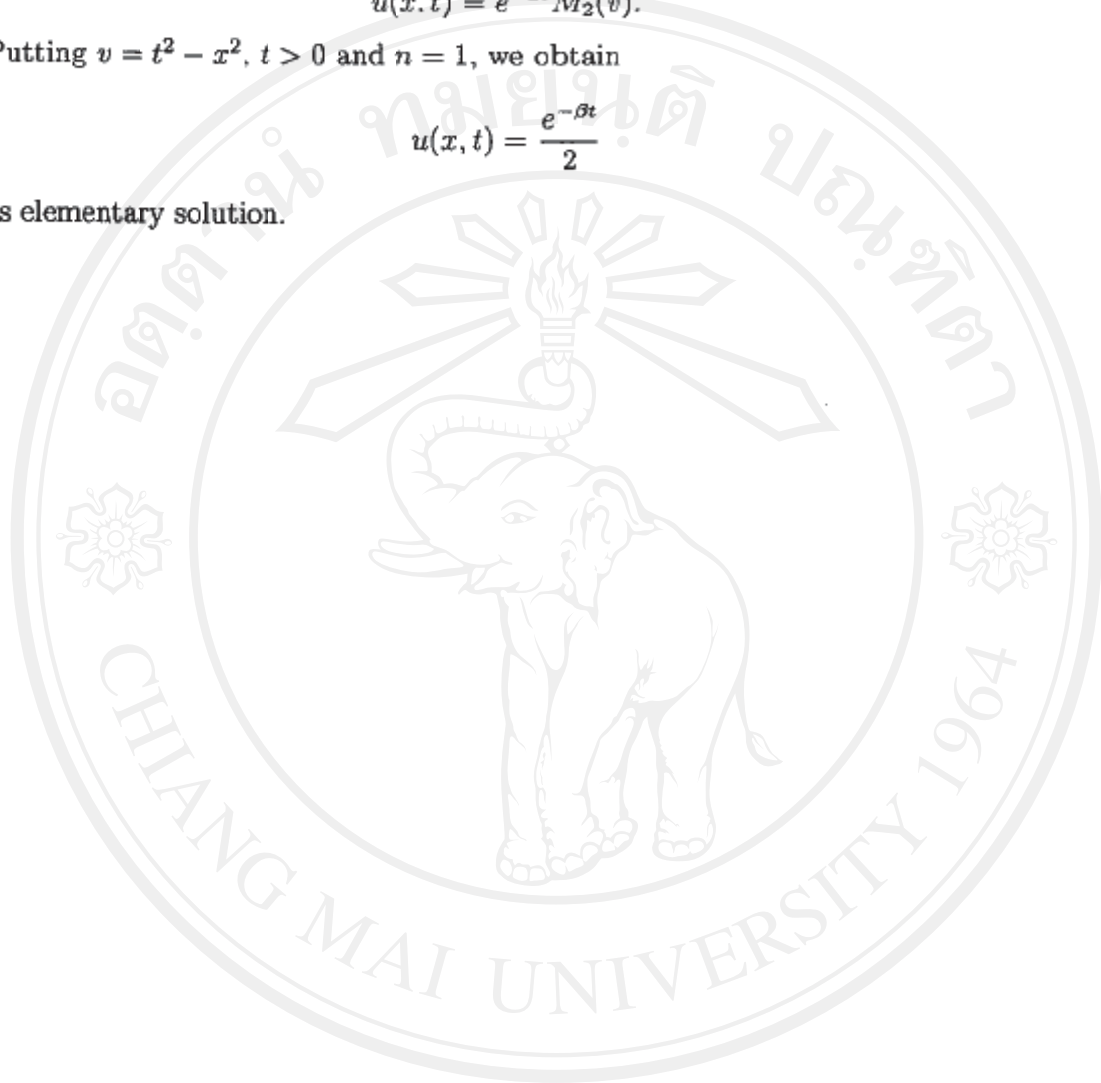
Now, we using the distribution  $e^{\alpha t} \square^k \delta$  for solve this problem. Recalling the solution of the telegraph equation (6.9) of the telegraph operator (6.1) with  $k = 1$ , we have

$$u(x, t) = e^{-\beta t} M_2(v).$$

Putting  $v = t^2 - x^2$ ,  $t > 0$  and  $n = 1$ , we obtain

$$u(x, t) = \frac{e^{-\beta t}}{2}$$

as elementary solution.



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